

# KINETIC AND QUANTUM PHENOMENA IN EXPANDING GAS FLOWS

V. V. Riabov

Rivier University, 420 South Main Street, Nashua, NH 03060, USA, vriabov@rivier.edu

This paper presents the results of numerical studies of rotational-translational nonequilibrium, kinetic, and diffusive processes in spherical expanding gas flows. Computations are made using the direct simulation Monte-Carlo (DSMC) method [1, 2] and the solver [3] of the Navier-Stokes equations in terms of classical and quantum concepts [3-7], at the Knudsen numbers  $Kn^*$  from 0.0015 to 0.03 and pressure ratios  $P = p_0/p_a$  from 100 to 10,000.

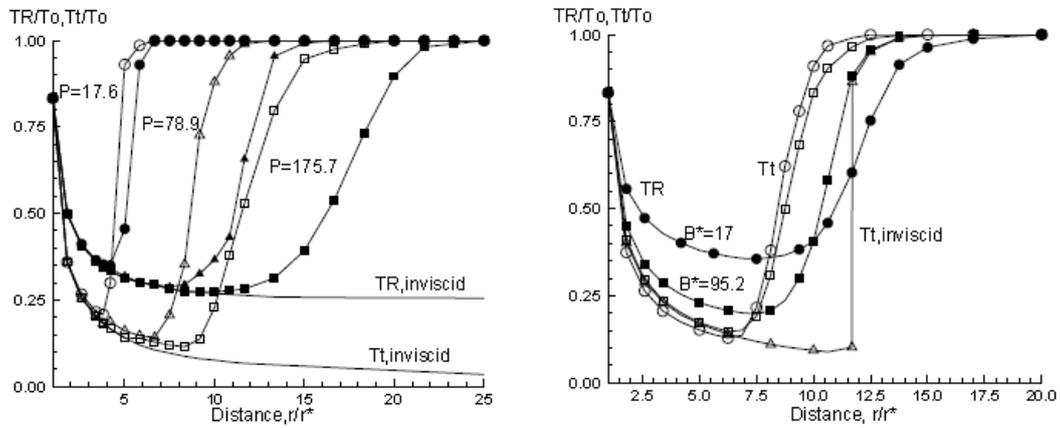


Fig. 1 *Left*: Rotational  $T_R$  (filled markers) and translational  $T_t$  (empty markers) temperatures in spherical flow at various pressure ratios:  $P = 17.6$  (circles),  $78.9$  (triangles), and  $175.7$  (squares). *Right*:  $T_R$  and  $T_t$  at various relaxation parameters:  $B^* = 17$  and  $95.2$ .

Experimental studies of underexpanded  $N_2$  jets [8, 9] discovered a delay of rotational temperature  $T_R$  compared to the translational one  $T_t$ . A drop in the gas density downstream leads to a decrease in the number of collisions and the  $T_R$  departure from the equilibrium value [6]. The Navier-Stokes equations and relaxation equation, based on the  $\tau$ -approximation [7], are solved by the numerical method [3]. Solutions depend on Reynolds number  $Re^*$ ,  $P = p_0/p_a$ , temperature  $T_a$ , and relaxation parameter  $B^*$  that is calculated as a ratio of  $p\tau_R$  and viscosity. Here index "a" refers to background conditions and index \* refers to sonic conditions. The distributions of  $T_R$  and  $T_t$  at  $Re^* = 161.83$ ,  $B^* = 28.4$ ,  $T_a = 295$  K, and various  $P$  values are shown in Fig. 1 (*left*). Computations confirm the delay of  $T_R$  compared to  $T_t$ . Inviscid flow parameters  $T_R$  and  $T_t$  were estimated in [6]. The decrease of the main relaxation parameter  $B^*$  leads to a faster "frozen" value of  $T_R$  in the supersonic-flow zone (see Fig. 1, *right*). The spherical flow could be separated by the coordinate  $R_S$ , at which the stream parameters are extreme, into two regions with different properties [10]. In the first "internal" region the flow is supersonic. The flow parameters depend on the Reynolds numbers  $Re^*$  and  $Re_S$  (or related Knudsen numbers  $Kn^*$  and  $Kn_S$ ) [6]. In the second "external" region, there is a transition of supersonic flow through the spherical shock wave into a subsonic stream. The Reynolds number  $Re_a$  (or related Knudsen number  $Kn_a$ ) based on the length scale parameter at infinity,  $R_a = r_*P^{1/2}$  [11], is the major similarity parameter in this region. The similarity factor  $K_2 = Re_a = Re^*/P^{1/2}$  [11] can be used to study the flow structure here. The major changes of  $T_R$  and  $T_t$  occur in the shock wave at values of the normalized distance parameter  $r/R_a$  about 1. The shock width decreases with increasing  $K_2$ .

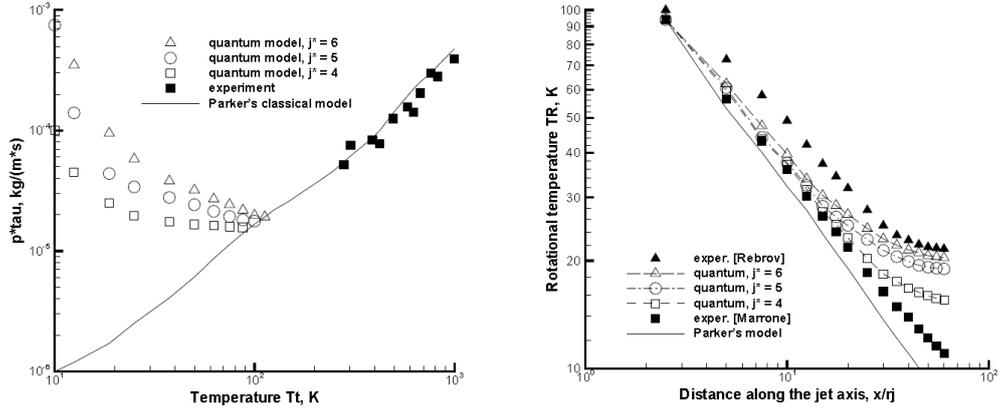


Fig. 2: *Left*: The rotational relaxation parameters  $p\tau_R(T_t)$  in  $N_2$  expanding into a vacuum. *Right*: The rotational  $T_R$  temperature in freejet expanding flow of  $N_2$ .

At the decrease of  $T_t$ , adiabatic collision conditions [12] should be taken into account, and the relaxation time  $p\tau_R$  increases due to the sharp decrease of the rotational transfer probabilities. Following [4],  $p\tau_R$  values were calculated for  $N_2$  at stagnation temperature  $T_0 = 295$  K [see Fig. 2, *left*] under the conditions of aerodynamic experiments in underexpanded jets [6-9]. At  $T_t > 273$  K, numerical results correlate well with experimental data [14, 15]. In the expansion of nitrogen, starting at  $T_0 = 300$  K, the maximum population of molecules occurs at rotational levels  $j^*$  from 6 to 4 [4]. The results of calculating  $p\tau_R$  for  $j^* = 6, 5$ , and 4 are shown in Fig. 2 (*left*). The calculations based on the classical concept [13] (see solid line in Fig. 2, *left*) do not show a tendency of increasing  $p\tau_R$  with the decrease of  $T_t$  at  $T_t < 100$  K.

For qualitative estimations, the energy relaxation time is replaced by the relaxation time of the level  $j^*$ . Figure 2 (*right*) shows the distributions of rotational temperature  $T_R$  along the axis of  $N_2$  jet at  $B_{*j} = (\rho u r / p\tau_R)_{*j} = 2730$ ,  $p_0 r_j = 240$  torr·mm and  $T_0 = 295$  K. The result of using the classical mechanics concept [13] (solid line) does not correlate the experimental data (filled squares [8], triangles [9]) for  $T_R$ , which are lower and upper bounds on the distribution of rotational energy along the  $N_2$  jet axis. Numerical results, based on the quantum concept for values of  $p\tau_R$  at  $j^* = 6$  and 5, correlate with the experimental data [9].

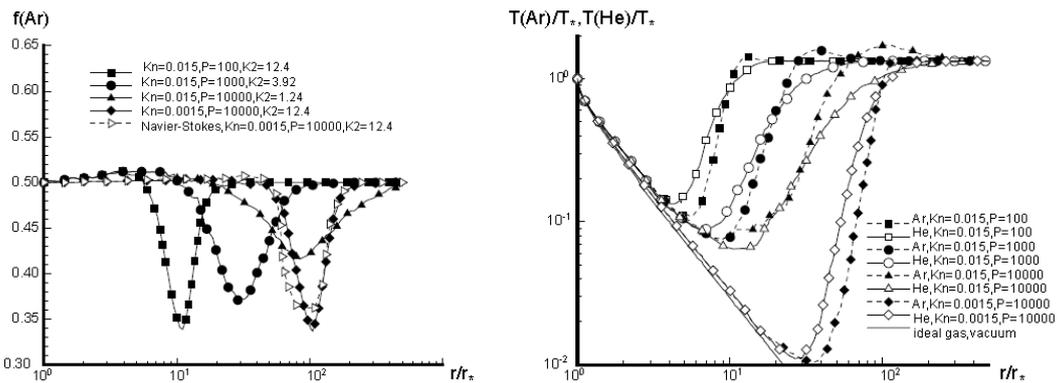


Fig. 3 Argon mole fraction distributions (*left*) and species temperatures (*right*) in spherical expanding flow of  $Ar-He$  mixture at different Knudsen numbers  $Kn^*$  and pressure ratios  $P$ .

Kinetic and diffusion effects in spherical expanding  $Ar-He$  flows (mole fraction  $f_{Ar,*} = 0.5$ ) were studied using the DSMC method [2] at  $Kn^*$  from 0.0015 ( $Re^* = 1240$ ) to 0.015 ( $Re^* = 124$ ) and  $P$  from 100 to 10,000. Both phenomena influence the shock thickness, parallel and transverse species' temperatures, diffusive velocities, and species separation. Distributions of  $f(Ar)$  and species temperatures are shown in Fig. 3. The species concent-

ration changes insignificantly at  $r < R_a$ . Accumulation of the light component occurs in the spherical shock (see Fig. 3, *left*) due to baro-diffusion effects [2]. For near-continuum flow conditions at  $K_2 = 12.4$ , the DSMC data correlates with solutions of the Navier–Stokes equations [16] (see Fig. 3, *left*). In contrast to the one-temperature continuum approach [16], the DSMC method allows simulating multi-temperature kinetic media. It is found that  $T(He)$  increases more rapidly than  $T(Ar)$  in the supersonic zone at  $r < R_a$  (see Fig3, *right*).

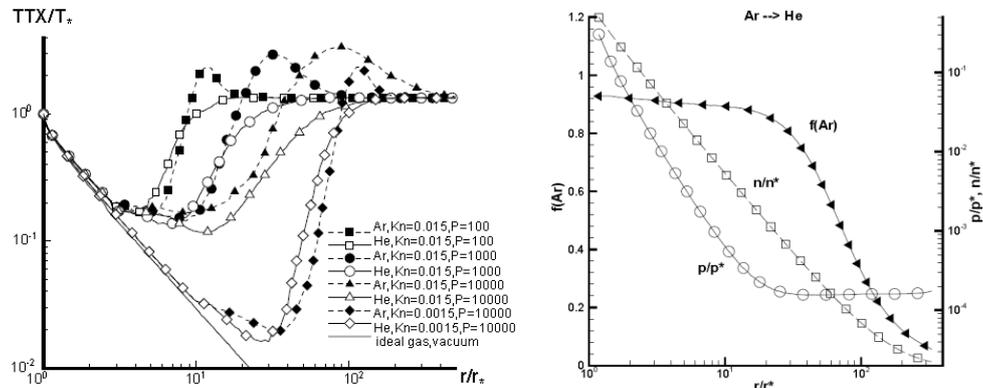


Fig. 4 *Left*: Parallel temperatures of species in spherical expanding flow of *Ar-He* mixture at different Knudsen numbers  $Kn_*$  and pressure ratios  $P$ . *Right*: Argon mole fraction, pressure and number density in expansion of argon into helium at  $Kn_* = 0.014$  and  $K_2 = 0.785$ .

In supersonic flow, the effect of freezing the parallel temperature  $TX$  found in [1] is confirmed (see Fig. 4, *left*). The freezing comes first for heavier molecules (*Ar*) at smaller values of  $Kn_*$ . The transverse temperature  $TY$  for both species follows the temperature in the isentropic expansion [1, 2]. The spherical expansion of a binary gas mixture into a flooded space was analyzed in the case of the presence of a diffusive flux at the infinity  $r \gg R_a$ . The numerical results were calculated for the case of the expansion of *Ar* with little *He* content ( $f_{Ar,*} = 0.99$ ) into a space filled by *He* with a small admixture of *Ar* ( $f_{Ar,a} = 0.02$ ). The distributions of argon concentration  $f_{Ar}$ , number density, and pressure at  $Kn_* = 0.014$ ,  $Re_* = 78.5$ , and  $K_2 = 0.785$  are shown in Fig. 4 (*right*). The case of the expansion of *He* into a space filled by *Ar* was analyzed in [2]. The results demonstrate that the background gas does not penetrate through the spherical shock wave into the inner supersonic region of the flow.

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